

Computational Efficiency of the Finite Element Method based on Second- Order Radiative Transfer Equation

Jun-Ming Zhao (赵军明)
J. Y. Tan and L. H. Liu

School of Energy Science and Engineering
Harbin Institute of Technology (HIT)

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Contents



Background

- Numerical insight into RTE
- Difficulties in numerical solving RTE
- Solution to the difficulties
- Second order RTE

Objective of this work

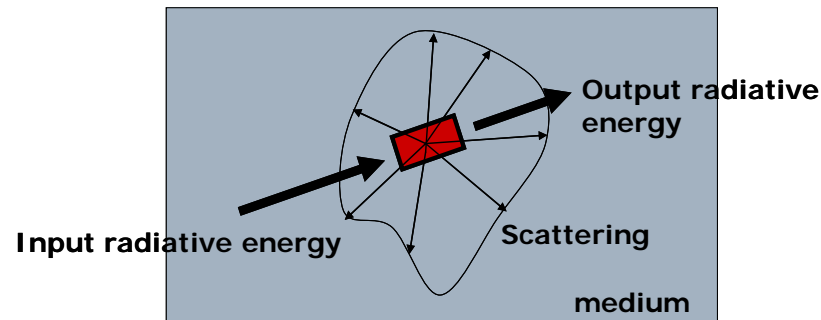
Formulation and Implementation

Results and Discussion

Numerical Insight into the RTE

First-Order Radiative Transfer Equation (FORTE, RTE) :

$$\boldsymbol{\Omega} \cdot \nabla I + (\kappa_a + \kappa_s) I = \kappa_a I_b + \frac{\kappa_s}{4\pi} \int_{4\pi} I(\mathbf{r}, \boldsymbol{\Omega}') \Phi(\boldsymbol{\Omega}, \boldsymbol{\Omega}') d\boldsymbol{\Omega}'$$



Boundary condition:

$$I(\mathbf{r}_w, \boldsymbol{\Omega}) = \bar{I}_{bw}(\mathbf{r}_w), \quad \mathbf{n}_w \cdot \boldsymbol{\Omega} > 0$$

Numerical Insight into the RTE



$$\boldsymbol{\Omega} \cdot \nabla I + \beta I = \kappa_a I_b + \frac{\kappa_s}{4\pi} \int_{4\pi} I(\mathbf{r}, \boldsymbol{\Omega}') \Phi(\boldsymbol{\Omega}', \boldsymbol{\Omega}) d\boldsymbol{\Omega}'$$

- If take direction as “velocity”, what it look like?

$$\mathbf{V} \cdot \nabla I + \beta I = S$$

Numerical Insight into the RTE



$$\boldsymbol{\Omega} \cdot \nabla I + \beta I = \kappa_a I_b + \frac{\kappa_s}{4\pi} \int_{4\pi} I(\mathbf{r}, \boldsymbol{\Omega}') \Phi(\boldsymbol{\Omega}', \boldsymbol{\Omega}) d\boldsymbol{\Omega}'$$

- If take direction as “velocity”, what it look like?

$$\mathbf{V} \cdot \nabla I + \alpha \nabla^2 I + \beta I = S$$



Numerical Insight into the RTE

$$\boldsymbol{\Omega} \cdot \nabla I + \beta I = \kappa_a I_b + \frac{\kappa_s}{4\pi} \int_{4\pi} I(\mathbf{r}, \boldsymbol{\Omega}') \Phi(\boldsymbol{\Omega}', \boldsymbol{\Omega}) d\Omega'$$

- If take direction as “velocity”, what it look like?

$$\mathbf{V} \cdot \nabla I + \alpha \nabla^2 I + \beta I = S$$

- *Due to the absence of diffusion term, it is a **Convection-dominated type** of general convection diffusion equation*

Difficulties in Numerical Solving RTE



1) *Stability problem*

- Caused by *Convection-dominated* characteristics of RTE



Difficulties in Numerical Solving RTE

1) *Stability problem*

- Caused by *Convection-dominated characteristics* of RTE

2) *False Scattering*

- Num. Phenomena: False energy diffusion
- Cause: Insufficient spatial accuracy



Difficulties in Numerical Solving RTE

1) *Stability problem*

- Caused by *Convection-dominated characteristics* of RTE

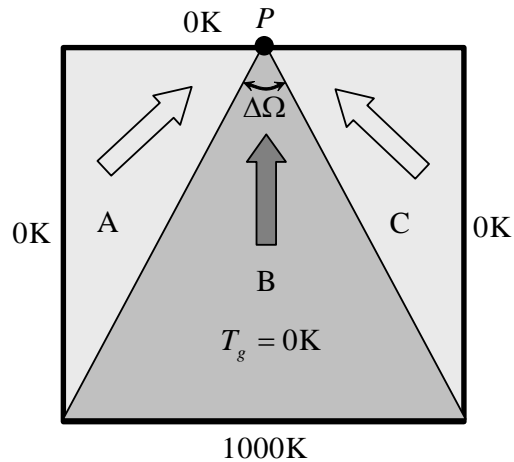
2) *False Scattering*

- Num. Phenomena: False energy scattering
- Cause: Insufficient spatial accuracy

3) “Ray Effects”

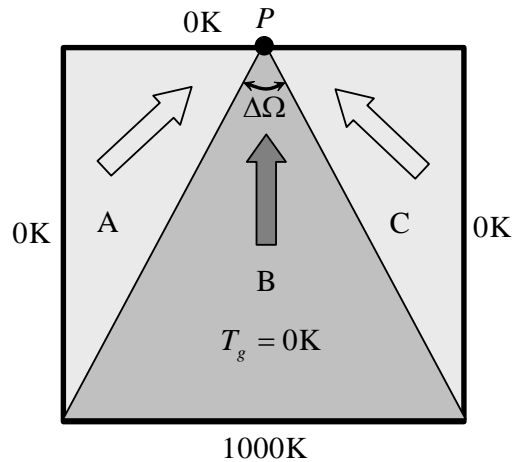
- Num. Phenomena: nonphysical wiggles in results
- Cause: Insufficient angular quadrature accuracy, **may coupled with 1) and 2)**

Example of Causes of *Ray effects*

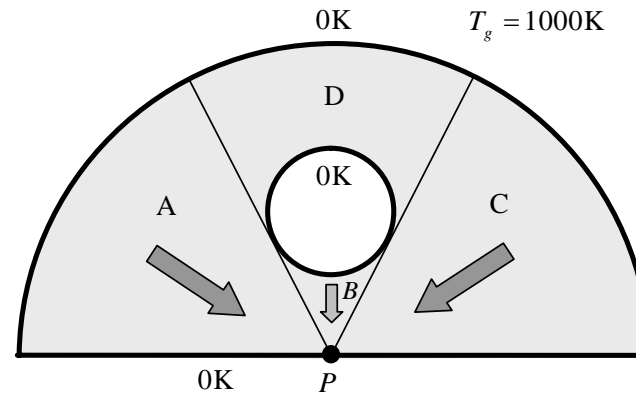


**Boundary load with
large nonuniformity**

Example of Causes of *Ray effects*

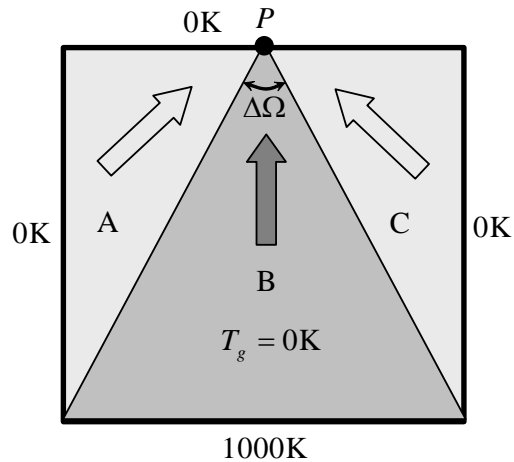


Boundary load with large nonuniformity

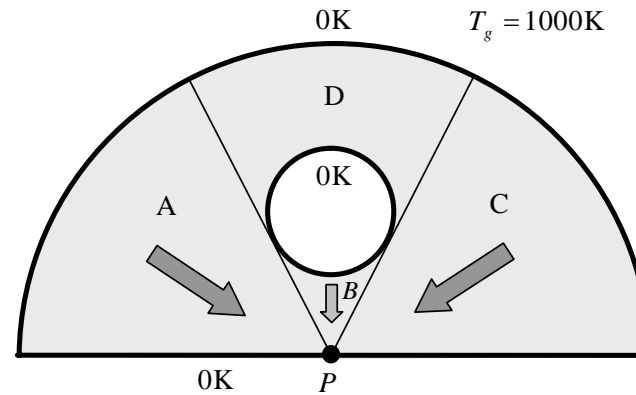


Interior Obstacle shielding

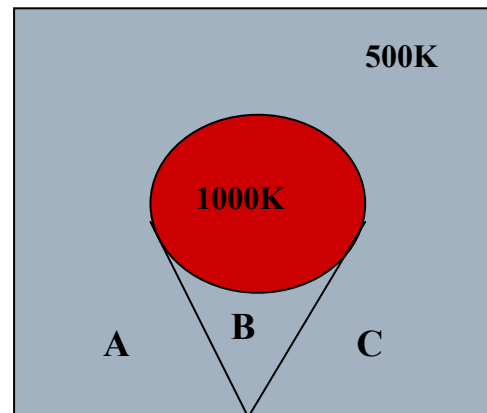
Example of Causes of *Ray effects*



Boundary load with large nonuniformity

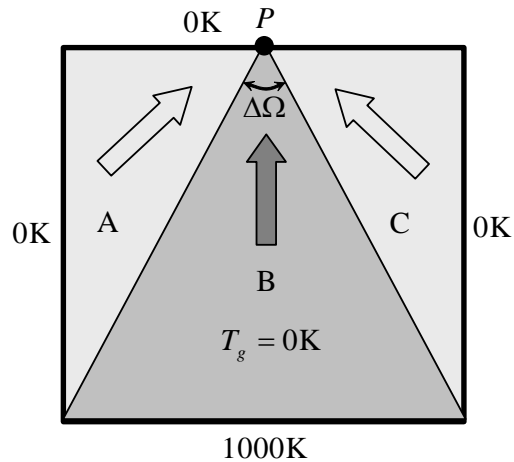


Interior Obstacle shielding

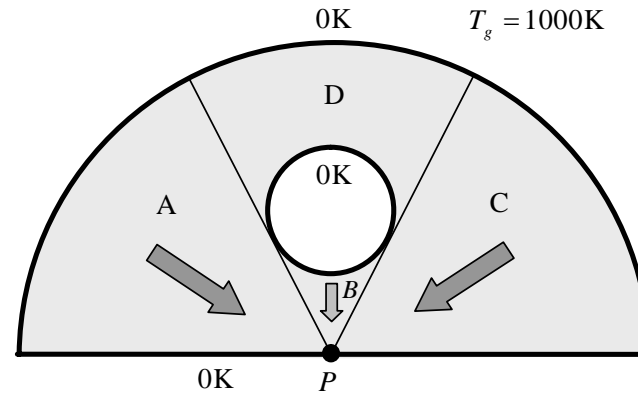


Large gradient of Source

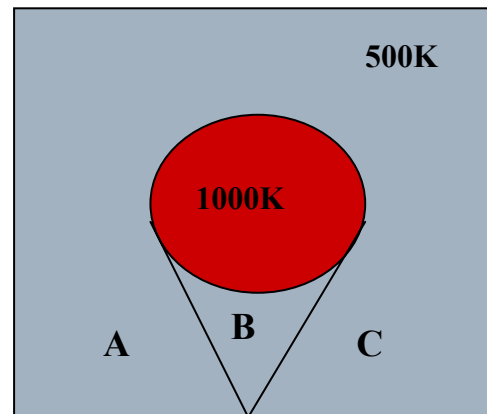
Example of Causes of *Ray effects*



Boundary load with large nonuniformity



Interior Obstacle shielding



Large gradient of Source

Zhao & Liu (JQSRT, 2007)



Solution to Difficulties

1) *Stability problem*

■ Basic rationale

- (A) Upwinding in discretization, artificial diffusion ...
- (B) Transform FORTE to second order equation, or cancel the convection term
 - Even-Parity formulation of RTE (EPF-RTE)
 - Second Order RTE (SORTE)

2) *False Scattering*

3) “Ray Effects”



Advantage of Solution (B)

- ❑ No artificial diffusion is needed to be intentionally added
- ❑ Based on the equation, radiative transfer can be solved stably with many methods, FEM, FVM, Meshless method,
- ❑ Hence it is a **unified approach**, one for all

Disadvantage of EPF-RTE



- ❑ Solution variable is not radiative intensity
- ❑ Difficult to be extended to anisotropic scattering media



Introduction to SORTE

□ Derivation

$$\frac{d}{ds} I + \beta I = \underbrace{\kappa_a I_b + \frac{\kappa_s}{4\pi} \int_{4\pi} I(\mathbf{r}, \Omega') \Phi(\Omega', \Omega) d\Omega'}_S$$

Introduction to SORTE



□ Derivation

$$\frac{1}{\beta} \frac{d}{ds} I + I = \frac{1}{\beta} S$$

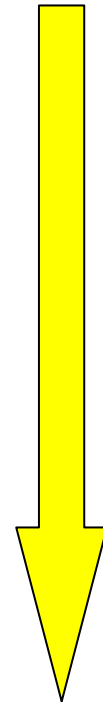
Introduction to SORTE

□ Derivation

$$\frac{d}{ds} \left[\frac{1}{\beta} \frac{d}{ds} I + I \right] = \frac{1}{\beta} S$$

$$\frac{d}{ds} \left[\frac{1}{\beta} \frac{d}{ds} I \right] + \frac{d}{ds} I = \frac{d}{ds} \left[\frac{1}{\beta} S \right]$$

$$-\frac{d}{ds} \left[\frac{1}{\beta} \frac{d}{ds} I \right] + \beta I = S - \frac{d}{ds} \left[\frac{1}{\beta} S \right]$$



Second Order Radiative Transfer Equation (SORTE) [Zhao & Liu(2007)]:

$$-\beta^{-1}\mathbf{\Omega}\cdot\nabla\left[\beta^{-1}\mathbf{\Omega}\cdot\nabla I\right]+I=S-\beta^{-1}\mathbf{\Omega}\cdot\nabla S$$

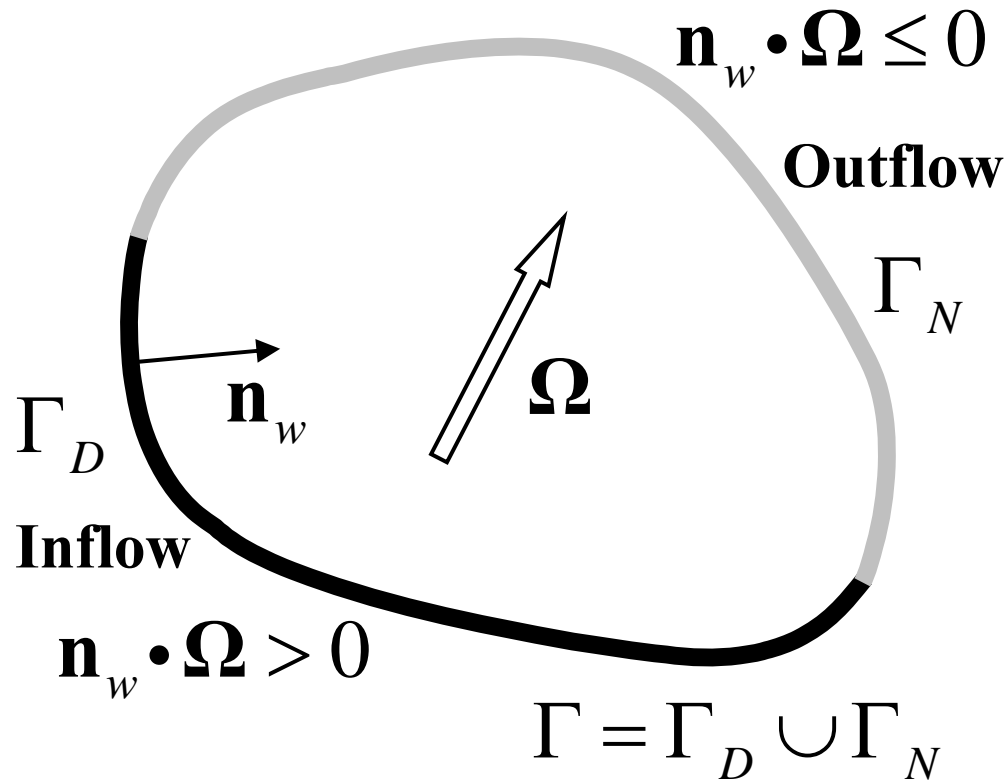
Properties of the SORTE

- Convection term is cancelled and replaced by a diffusion term
- Solution variable is intensity
- Can be easily applied to anisotropic scattering media, without limit on general applicability of FORTE

Introduction to SORTE



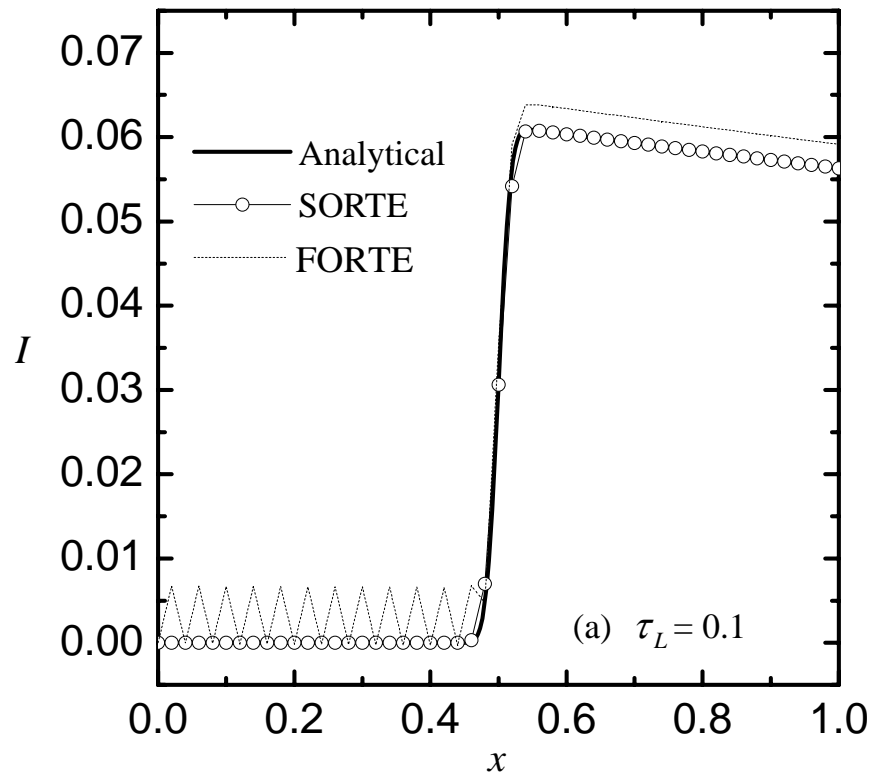
Boundary conditions



SORTE Performance Demo.



Solved intensity distribution in a slab with a Gaussian hill source



[Zhao & Liu(2007)]



Then, However,

- *What is the weakness of the **SORTE**?*
 - *Probably the most important is the **computational efficiency** of the numerical methods based on it*
- *Which is crucial for broad application of this approach.*
- *As such, this subject forms the major motivation of present research*



Objectives of this work

Investigate the

accuracy and **solution cost** of finite element method (FEM) based on the **SORTE**



Formulation and Implementation

SORTE in 2D can be written as:

$$(\mu^m)^2 \frac{\partial^2 I^m}{\partial x^2} + (\eta^m)^2 \frac{\partial^2 I^m}{\partial y^2} + 2\mu^m \eta^m \frac{\partial^2 I^m}{\partial x \partial y} - \beta^2 I^m = U^m$$

$$U^m = \mathbf{\Omega}^m \cdot \nabla S^m - \beta S^m$$

Formulation and Implementation

FEM discretization of the SORTE:

- 1) **FEM** approximation

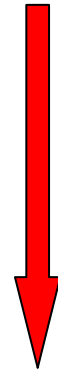
$$I^m(\mathbf{r}) \approx \sum_{i=1}^{N_{sol}} I_i^m \phi_i(\mathbf{r})$$

- 2) **Galerkin** approach

$$\sum_{i=1}^{N_{sol}} I_i^m \int_V \left[(\mu^m)^2 \frac{\partial^2 I^m}{\partial x^2} + (\eta^m)^2 \frac{\partial^2 I^m}{\partial y^2} + 2\mu^m \eta^m \frac{\partial^2 I^m}{\partial x \partial y} - \beta^2 I^m \right] \phi_j(\mathbf{r}) dV$$

$$= \int_V U^m(\mathbf{r}) \phi_j(\mathbf{r}) dV$$

- 3) final matrix form: $\mathbf{K}^m \mathbf{I}^m = \mathbf{H}^m$





Formulation and Implementation

By using tool matrices approach [Zhao & Liu (2006)]:

$$\mathbf{K}^m = (\mu^m)^2 \mathbf{A}^{xx} + \mu^m \eta^m \mathbf{A}^{xy} + \eta^m \mu^m (\mathbf{A}^{xy})^T + (\eta^m)^2 \mathbf{A}^{yy} + \beta^2 \mathbf{B}^{oo} + \beta (\mu^m \mathbf{N}^x + \eta^m \mathbf{N}^y)$$

$$\mathbf{H}^m = \left[\beta \mathbf{B}^{oo} - \mu^m (\mathbf{B}^{xo})^T - \eta^m (\mathbf{B}^{yo})^T + \beta (\mu^m \mathbf{N}^x + \eta^m \mathbf{N}^y) \right] \mathbf{S}^m$$

$$A_{ji}^{xx} = \int_V \frac{\partial \phi_j}{\partial x} \frac{\partial \phi_i}{\partial x} dV \quad A_{ji}^{xy} = \int_V \frac{\partial \phi_j}{\partial x} \frac{\partial \phi_i}{\partial y} dV \quad A_{jn}^{yy} = \int_V \frac{\partial \phi_j}{\partial y} \frac{\partial \phi_i}{\partial y} dV$$

$$B_{ji}^{xo} = \int_V \frac{\partial \phi_j}{\partial x} \phi_i dV \quad B_{jn}^{yo} = \int_V \frac{\partial \phi_j}{\partial y} \phi_i dV \quad B_{jn}^{oo} = \int_V \phi_j \phi_i dV$$

$$N_{ji}^x = \int_{\Gamma_N} \phi_j \phi_i (\mathbf{n}_w \cdot \mathbf{i}) dA \quad N_{ji}^y = \int_{\Gamma_N} \phi_j \phi_i (\mathbf{n}_w \cdot \mathbf{j}) dA$$

Formulation and Implementation

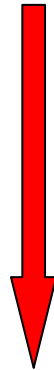
FEM discretization of the FORTE:

- 1) *FEM* approximation

$$I^m(\mathbf{r}) \approx \sum_{i=1}^{N_{sol}} I_i^m \phi_i(\mathbf{r})$$

- 2) *Galerkin* and *Least Squares* approach
- 3) final matrix form:

$$\mathbf{K}^m \mathbf{I}^m = \mathbf{H}^m$$



Formulation and Implementation



FORTE with Galerkin scheme:

$$\mathbf{K}^m = \mu^m (\mathbf{B}^{xo})^T + \eta^m (\mathbf{B}^{yo})^T + \beta \mathbf{B}^{oo}$$

$$\mathbf{H}^m = \mathbf{B}^{oo} \mathbf{S}^m$$

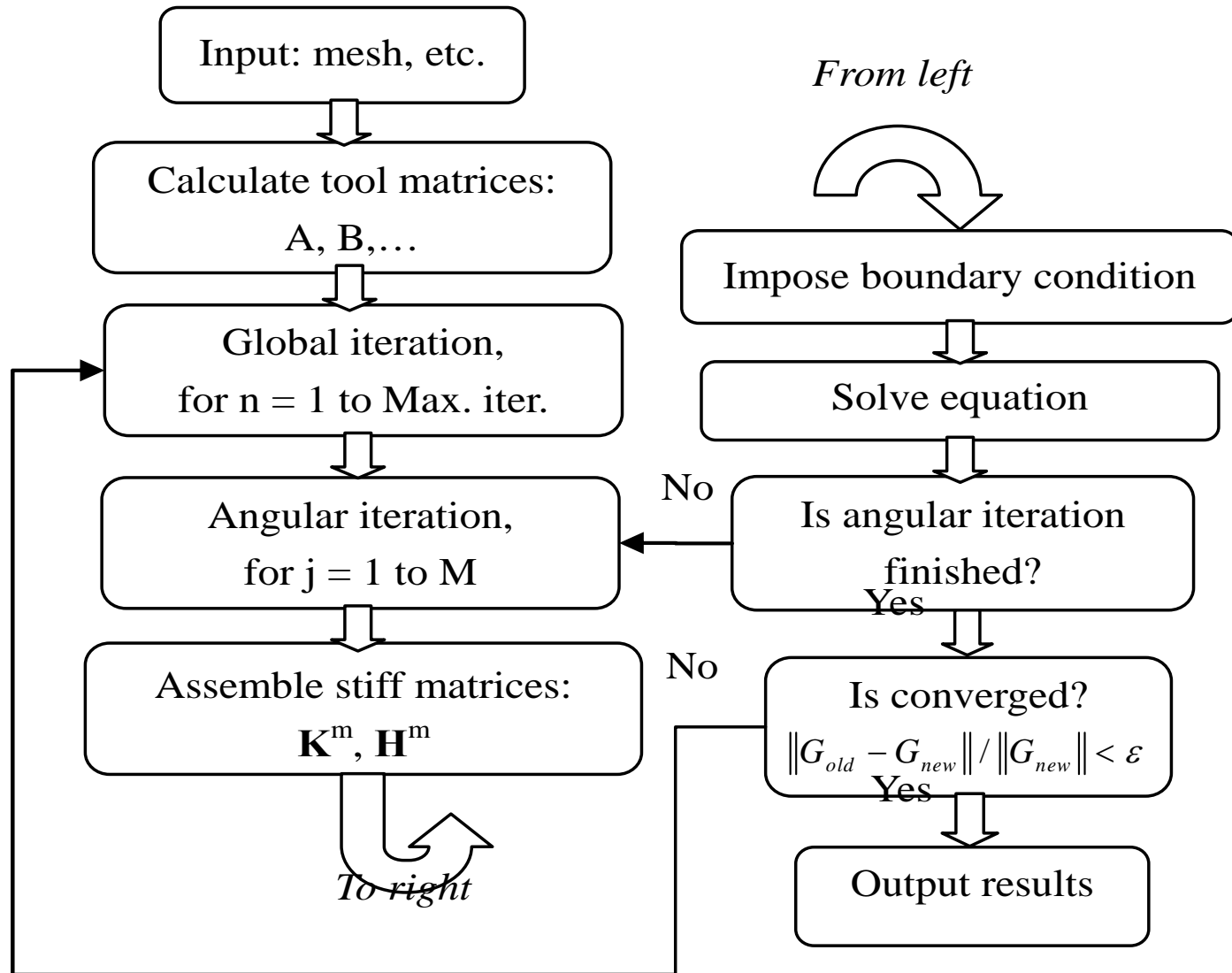
Formulation and Implementation

FORTE with Least square scheme:

$$\begin{aligned} \mathbf{K}^m = & (\mu^m)^2 \mathbf{A}^{xx} + \mu^m \eta^m \mathbf{A}^{xy} + \mu^m \beta \mathbf{B}^{xo} \\ & + \eta^m \mu^m (\mathbf{A}^{xy})^T + (\eta^m)^2 \mathbf{A}^{yy} + \eta^m \beta \mathbf{B}^{yo} \\ & + \beta \mu^m (\mathbf{B}^{xo})^T + \beta \eta^m (\mathbf{B}^{yo})^T + (\beta)^2 \mathbf{B}^{oo} \end{aligned}$$

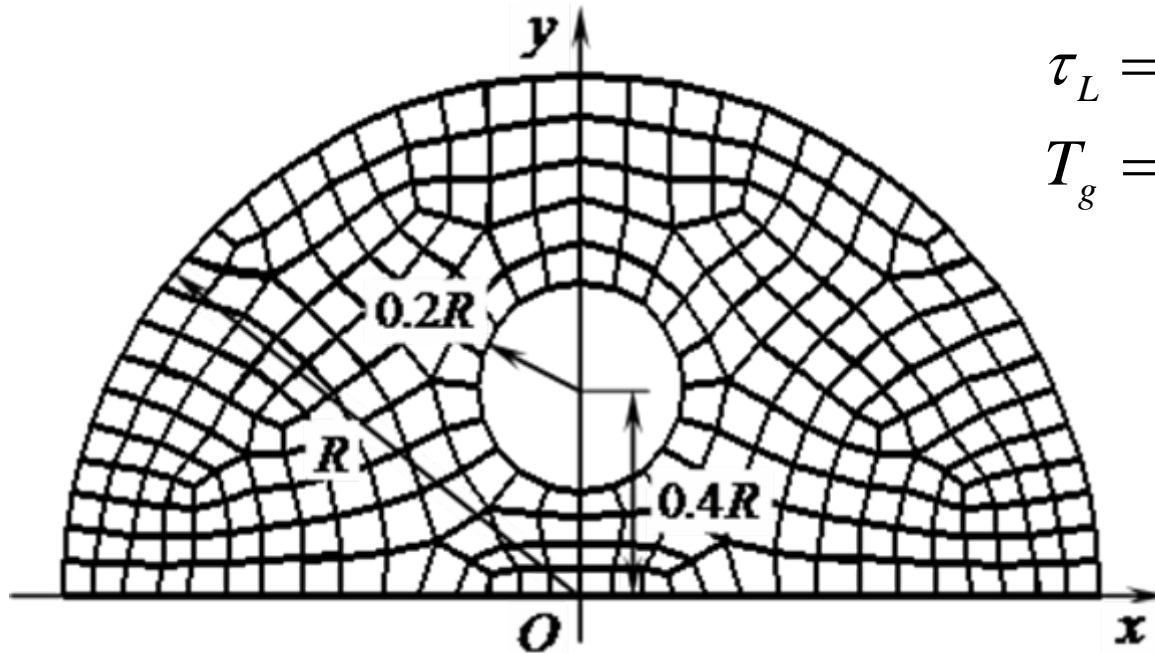
$$\mathbf{H}^m = \left(\mu^m \mathbf{B}^{xo} + \eta^m \mathbf{B}^{yo} + \xi^m \mathbf{B}^{zo} + \beta \mathbf{B}^{oo} \right) \mathbf{S}^m$$

Generic Solution procedures



Results and Discussion

Case 1: Semicircular enclosure with a circular hole

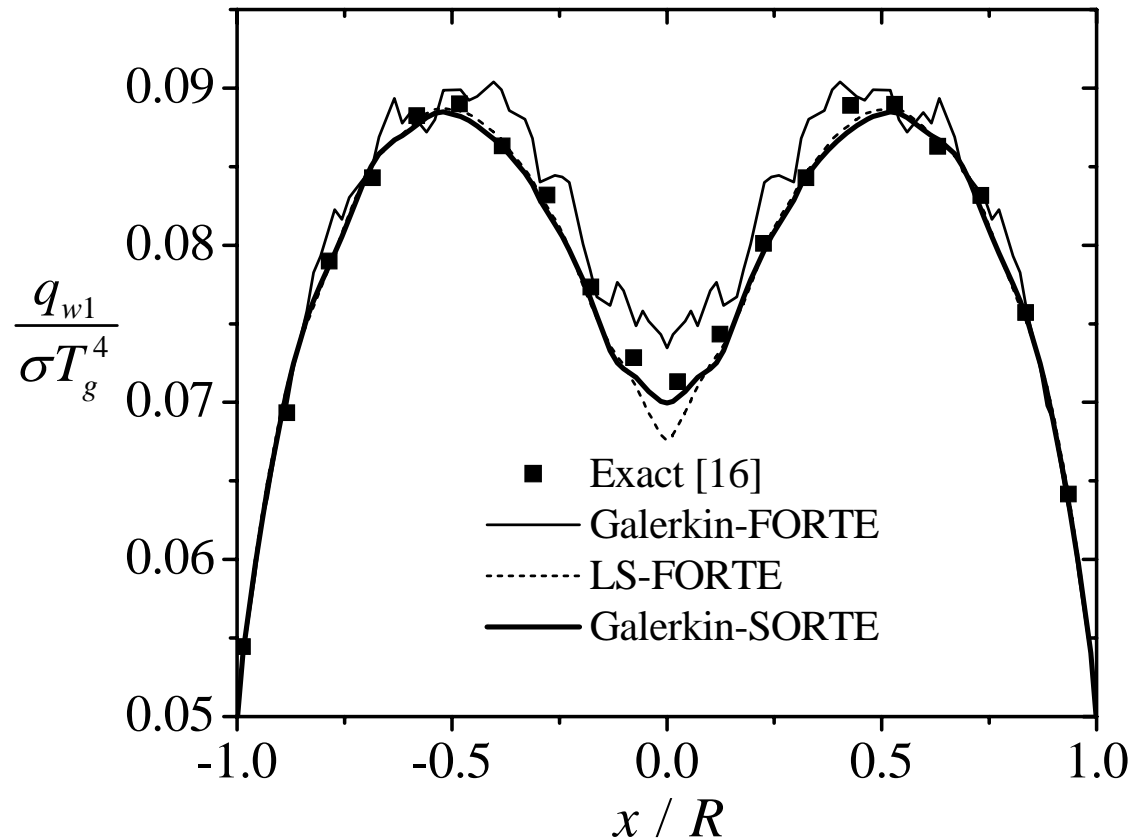


$$\tau_L = \beta R = 0.1$$

$$T_g = 1000 \text{ K}$$

Configuration of the semicircular enclosure and mesh decomposition (272 elements).

Results and Discussion



Heat flux distribution along bottom wall

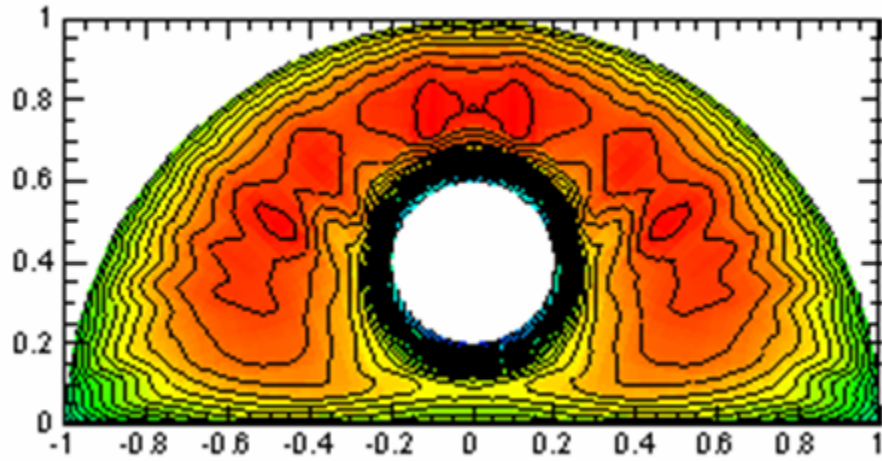
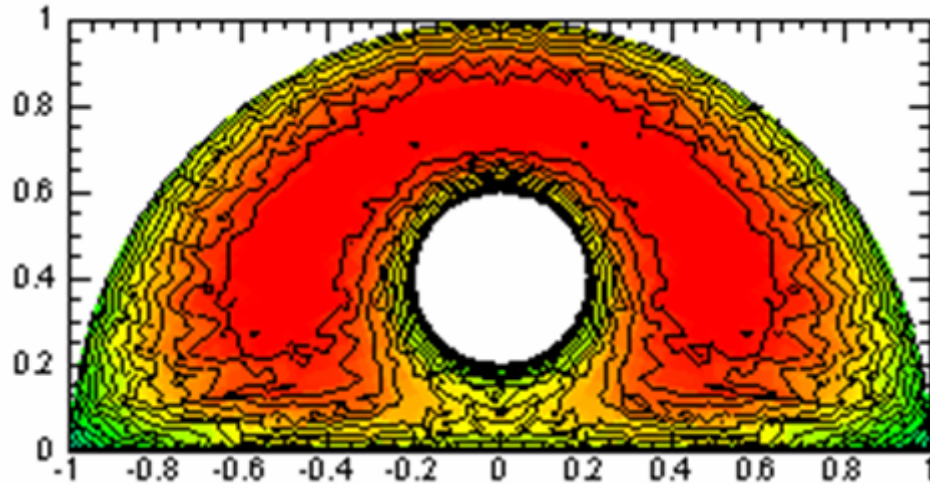
Space: 272 elements,
shape function is
constructed
through 3rd order
Chebyshev
approximation,

Solid angle:

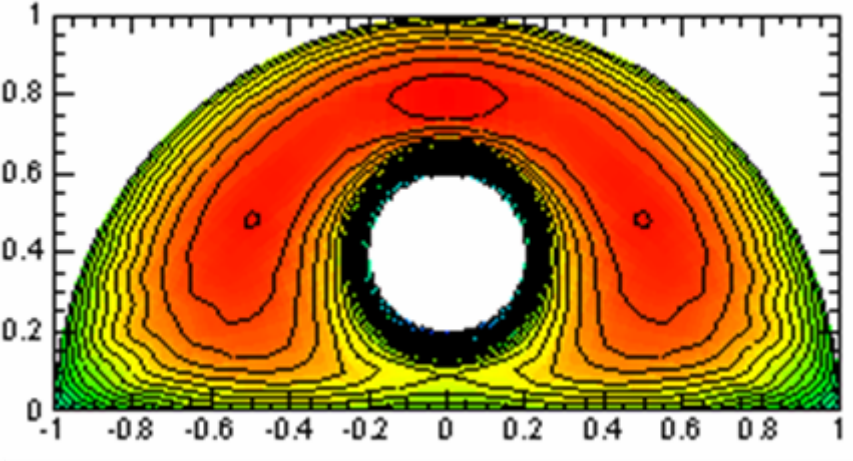
$$N_{\theta} \times N_{\phi} = 20 \\ \times 40$$

Results and Discussion

**Galerkin-
FORTE**



LS-FORTE

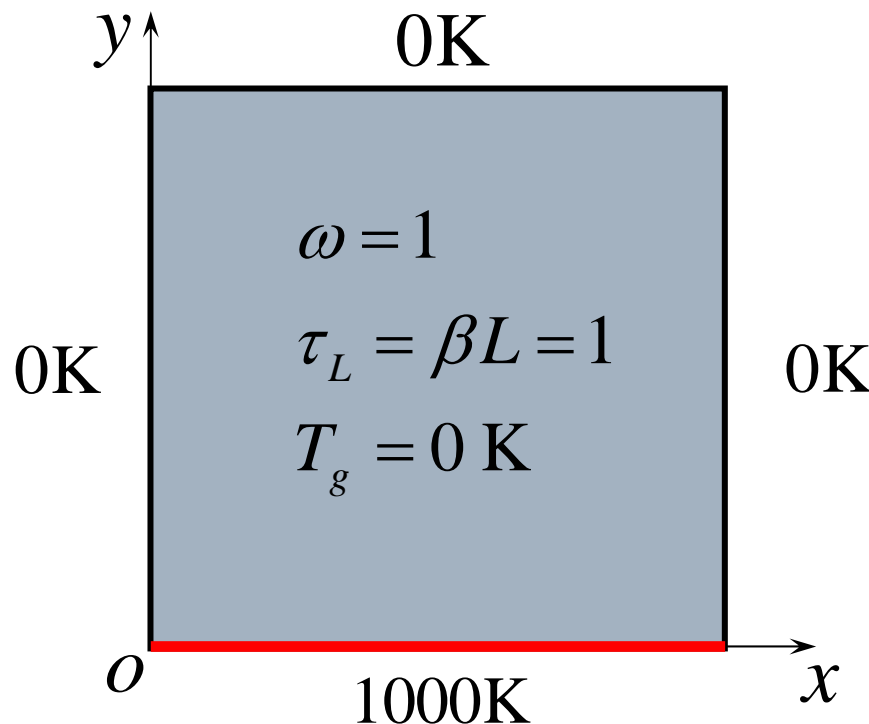


Galerkin-SORTE



Results and Discussion

Case 2: Isotropically Scattering Medium in a Square Enclosure



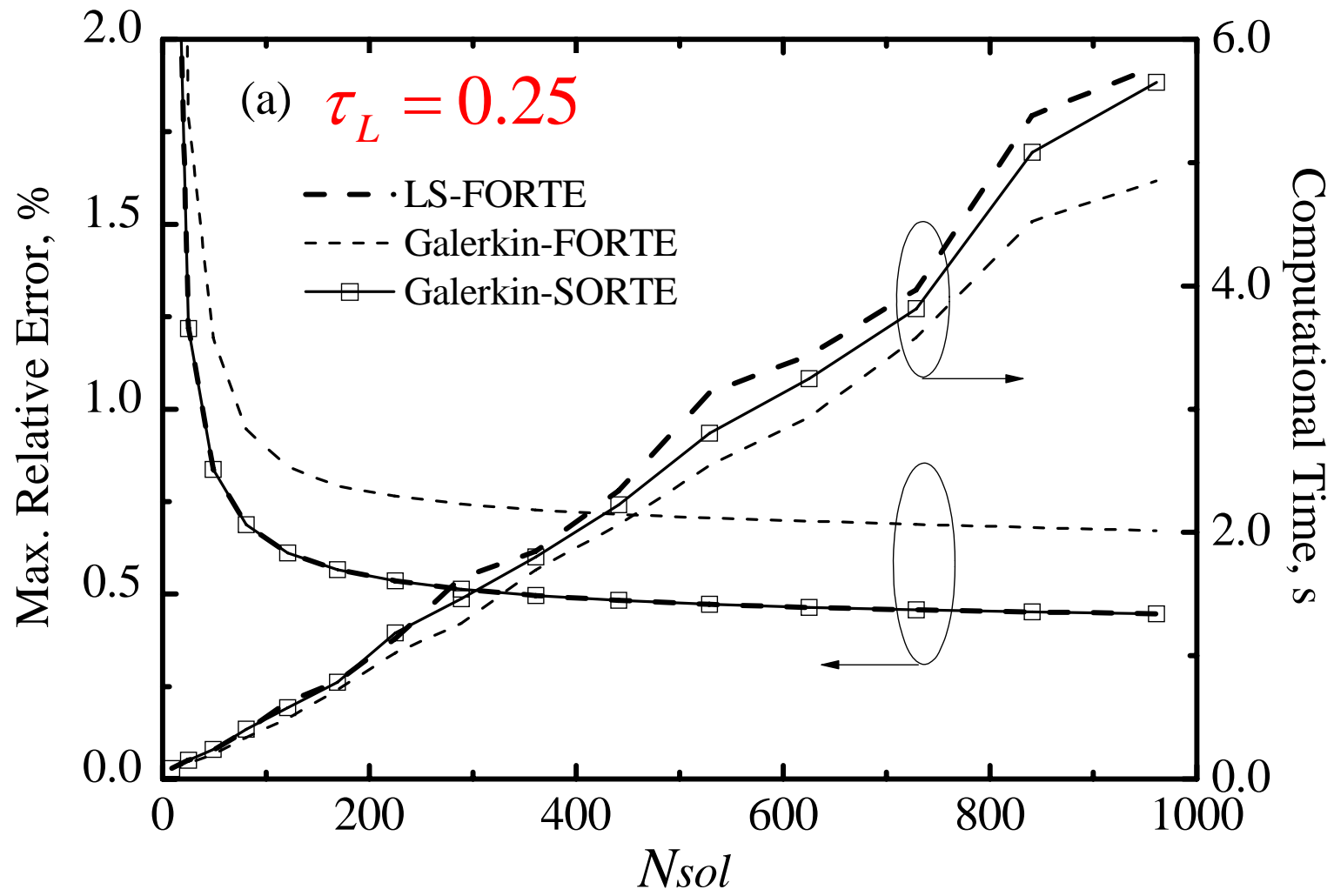
Space: $M \times M$

bilinear elements,
 M is taken as
needed for
convergence
analysis

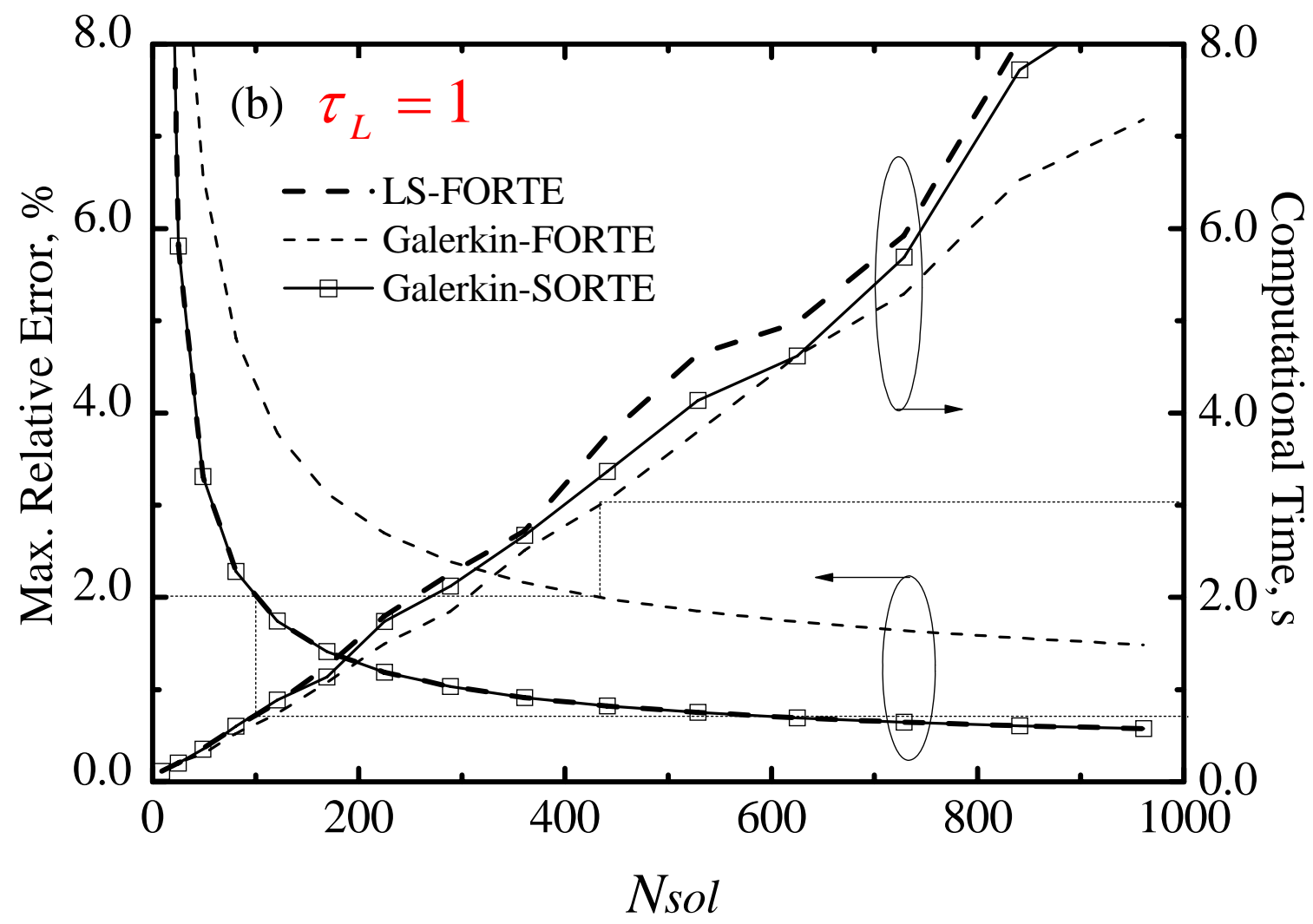
Solid angle: S_8

Solution quantity:
bottom wall
radiative heat flux

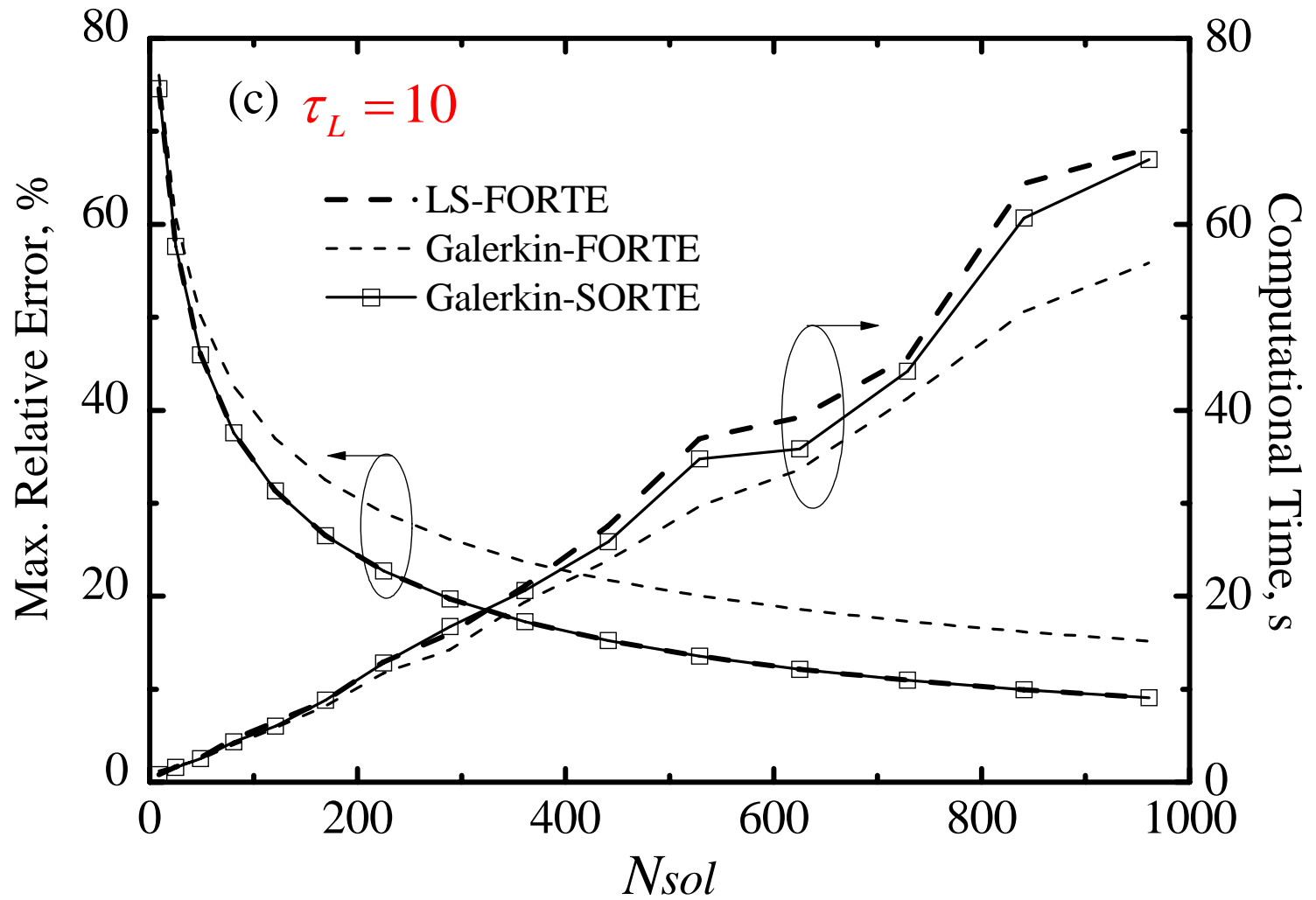
Results and Discussion



Results and Discussion



Results and Discussion



Conclusions



- The **accuracy** of the **FEM** based on the **SORTE** is generally better than that based on the FORTE
- *FEM based on SORTE* is the *most efficient* than the FEMs based on the FORTE.

Questions & comments?

*Thanks for your
attention!*

